

Memory-efficient nonequilibrium Green's function framework built on quantics tensor trains

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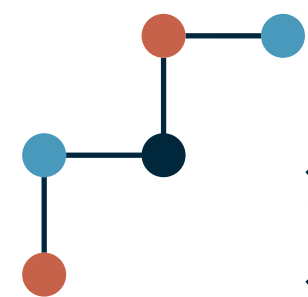
Anna Kauch



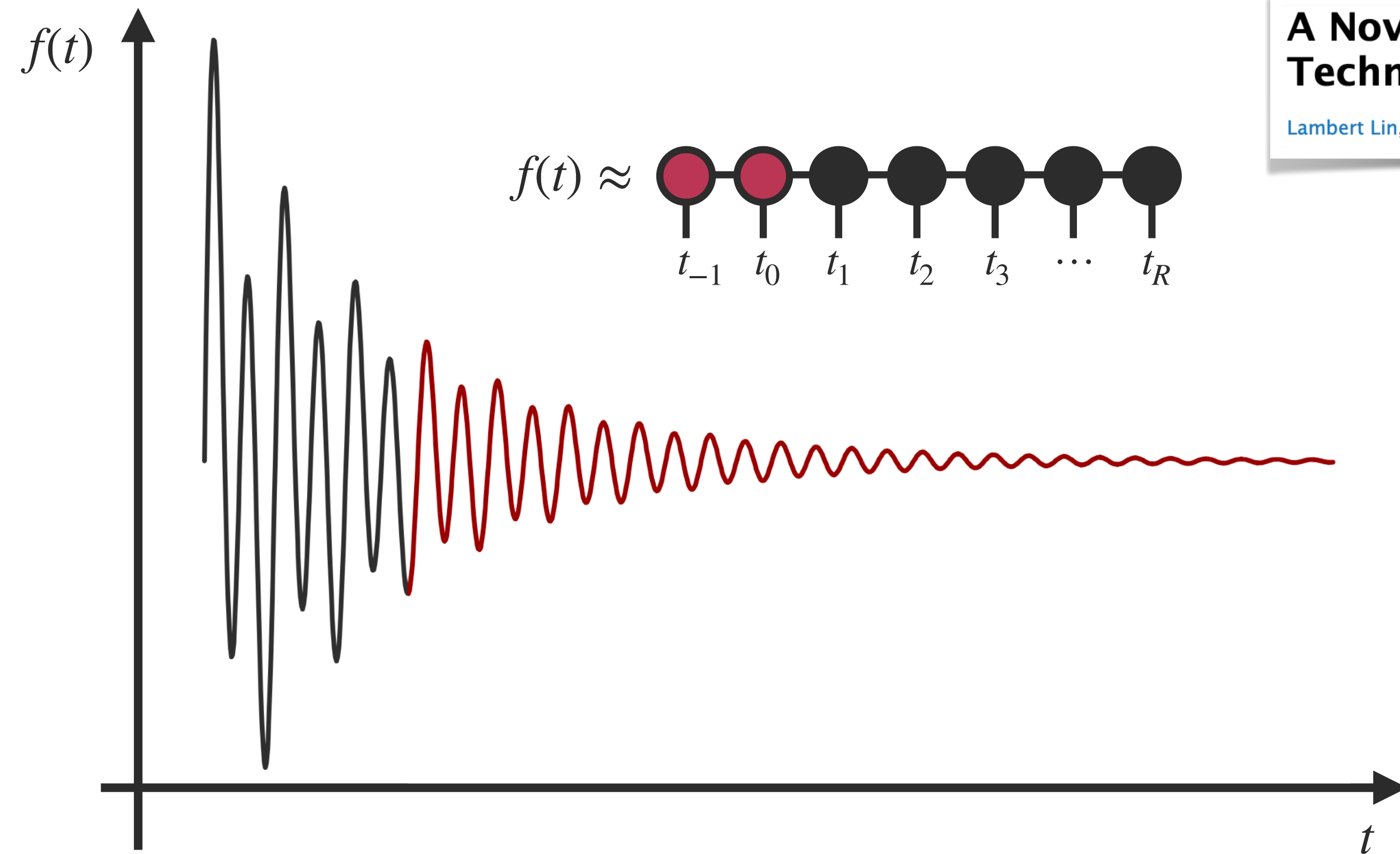
Hiroshi Shinaoka



Philipp Werner



**Swiss National
Science Foundation**



A Novel Method of Function Extrapolation Inspired by Techniques in Low-entangled Many-body Physics

Lambert Lin, Steven R White

...?

dynamic mode decomposition

$$\frac{d\mathbf{x}(t)}{dt} = \overset{\text{nonlinear}}{\mathbf{f}(\mathbf{x}(t), t)} \approx \overset{\text{linear}}{\mathbf{A}\mathbf{x}(t)}$$

Schmid, J. Fluid Mech. 656, 5–28 (2010)

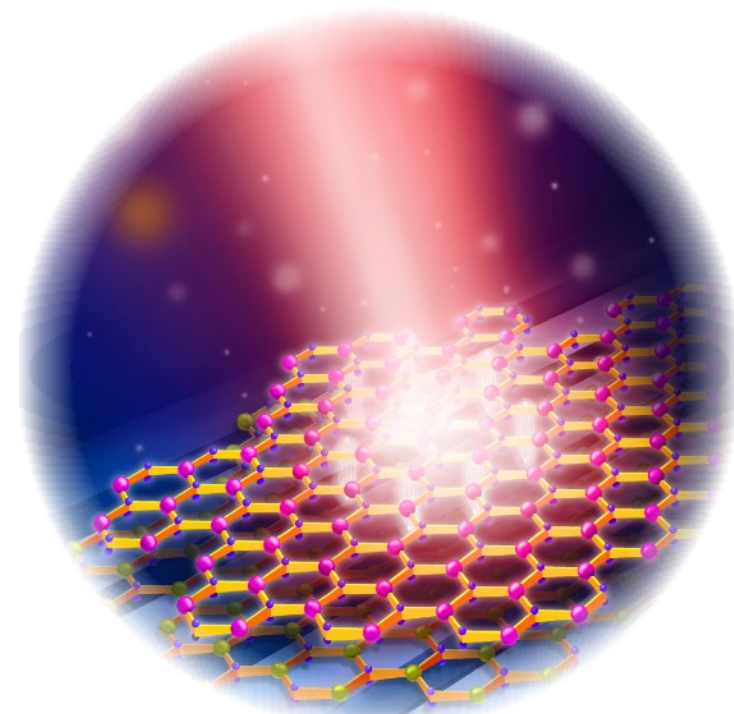
low-rank

$$\mathbf{x}(t) \approx \sum_{\ell} \phi_{\ell} \exp(i\omega_{\ell} t) b_{\ell}$$

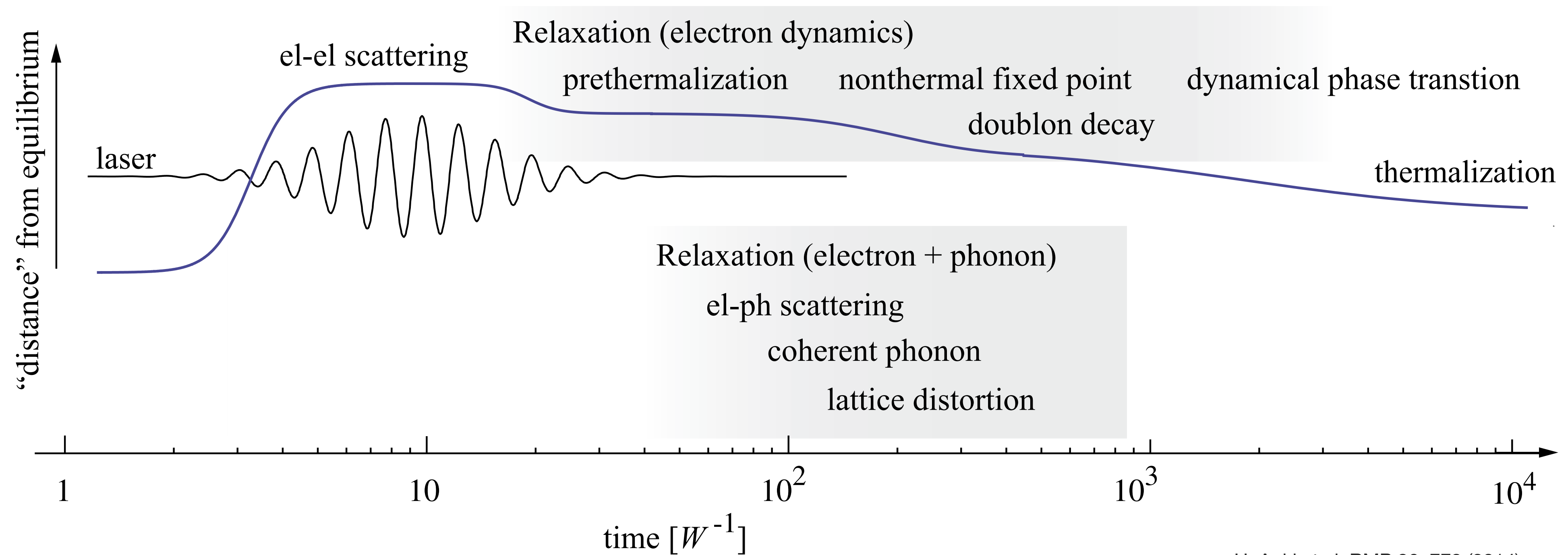
eigenvectors and eigenvalues of $\mathbf{A}$

DMD integrates into the QTT framework, maintaining exponentially fine resolution, and offering **extrapolation, interpolation, and denoising**

M. Środa et al, arXiv:2509.22177, 2025

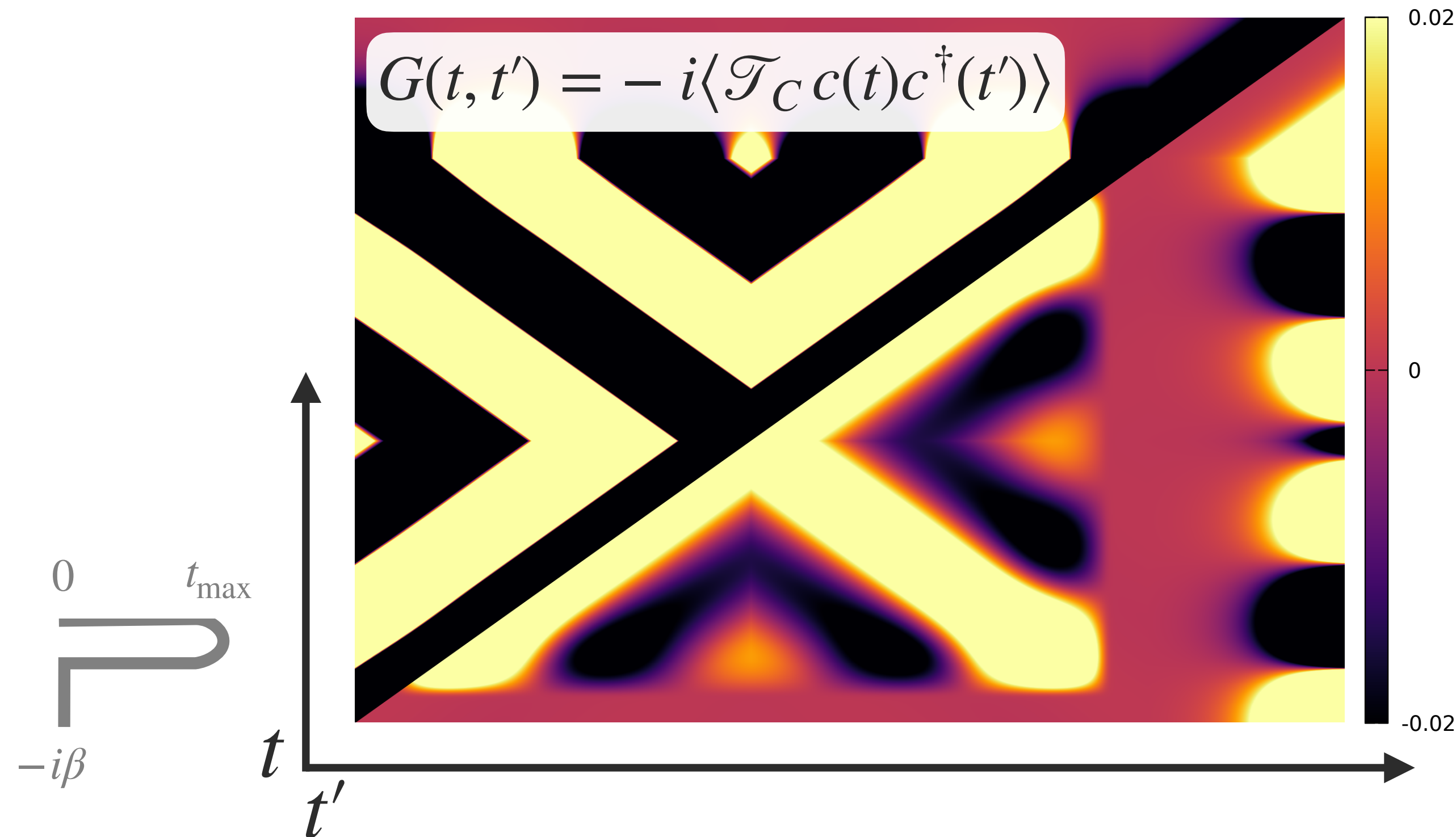


Larrea Young/larreayoung.com



H. Aoki et al, RMP 86, 779 (2014)

Nonequilibrium Green's function (NEGF) formalism faces a memory bottleneck



Access to:

- ✓ single-particle expectation values
- ✓ photoemission spectrum
- ✓ energy

Dyson equation

$$G = G_0 + G_0 \star \Sigma \star G$$

$$\int_C d\bar{t} \Sigma(t, \bar{t}) G(\bar{t}, t') \simeq \sum_{\bar{t}} \Delta \bar{t} \Sigma_{t, \bar{t}} G_{\bar{t}, t'}$$

convolution = matrix multiplication
 need to keep all previous times in memory

Ongoing work to avoid this

SciPost

SciPost Phys. 10, 091 (2021)

Low rank compression in the numerical solution of the nonequilibrium Dyson equation

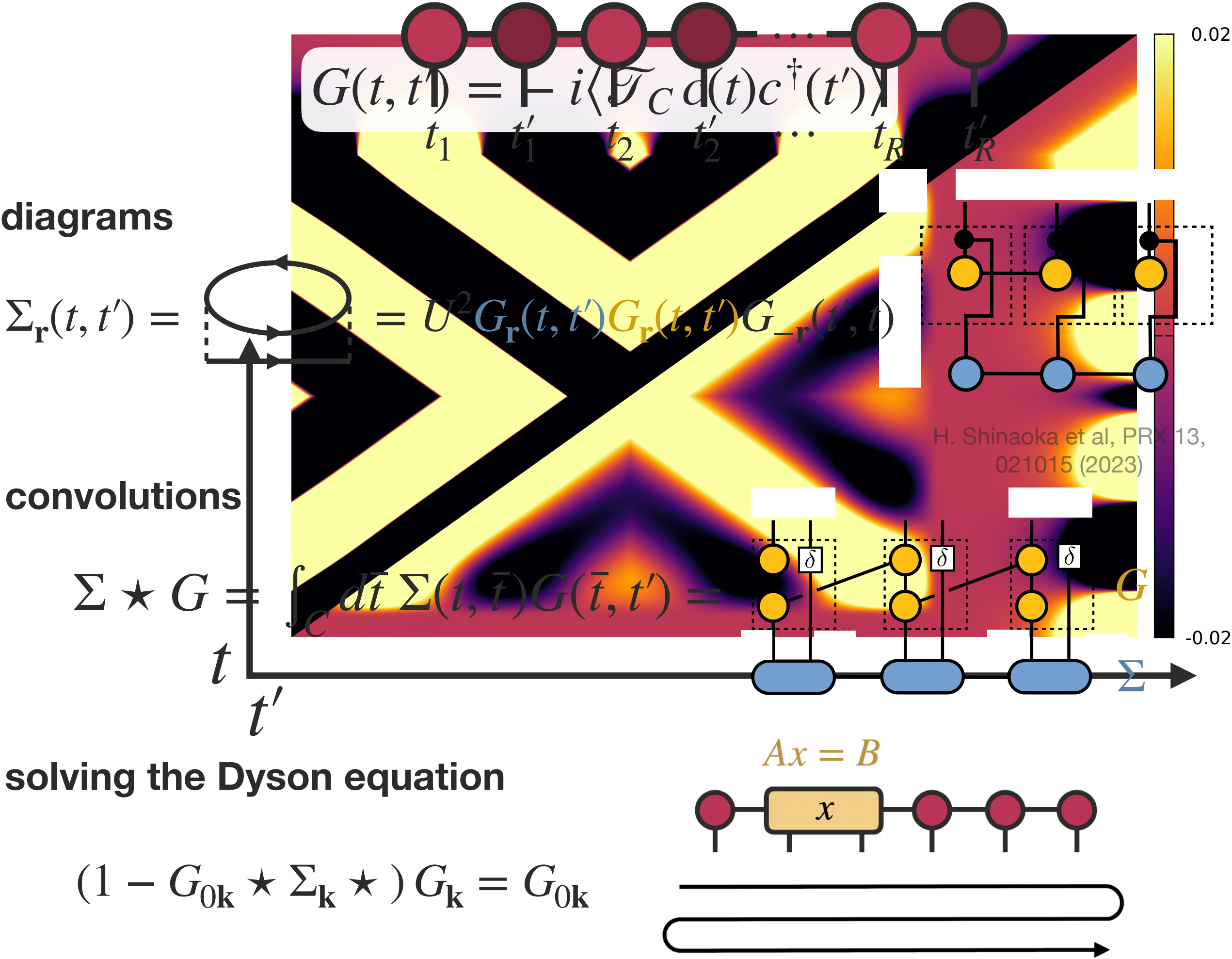
Jason Kaye^{1,2} and Denis Golež^{2,3,4}

Memory truncated Kadanoff-Baym equations

Christopher Stahl, Nagamalleswararao Dasari, Jiajun Li, Antonio Picano, Philipp Werner, and Martin Eckstein
 Phys. Rev. B **105**, 115146 – Published 31 March 2022

■ ■ ■

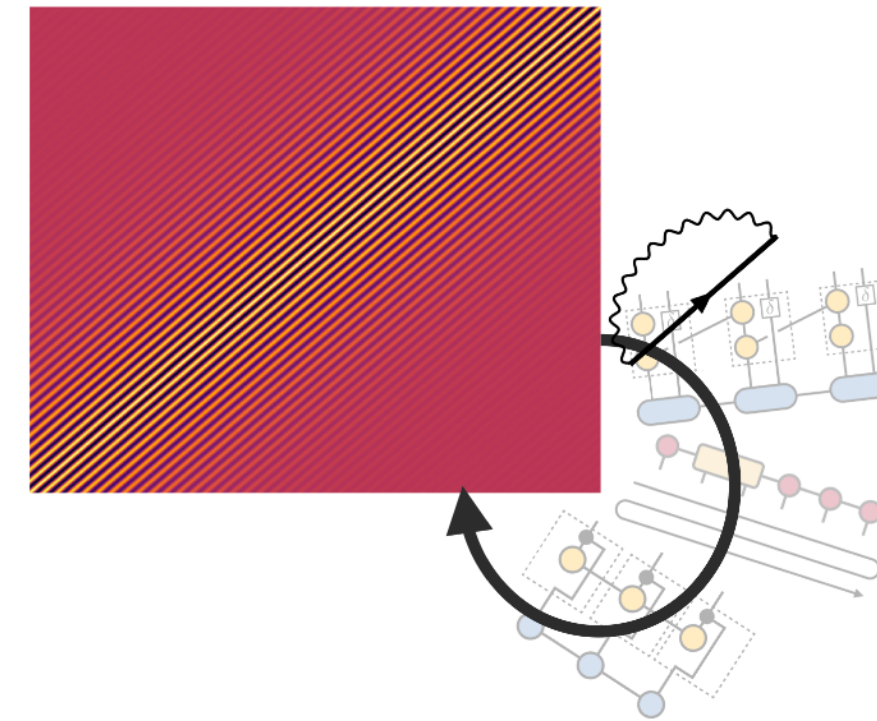
Quantics tensor trains outcompete conventional matrix-based approaches, but the convergence is not satisfactory



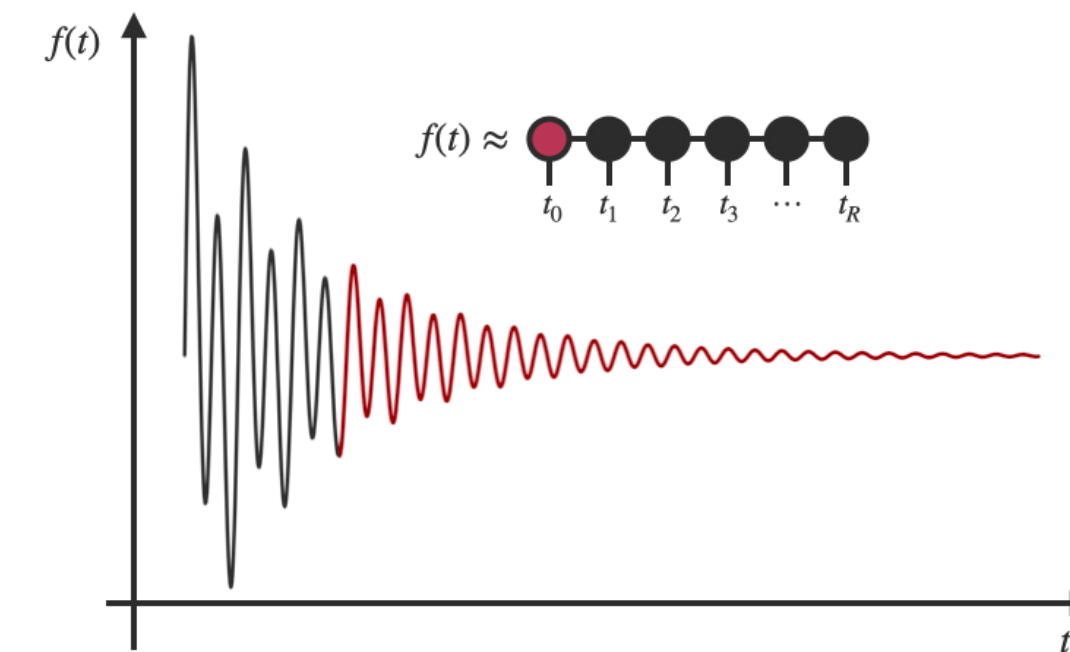
PROBLEM: unstable and slow convergence, difficulty extending $t_{\max}$ and increasing U

SOLUTION: reliable extrapolated initial guess and causality-based solver

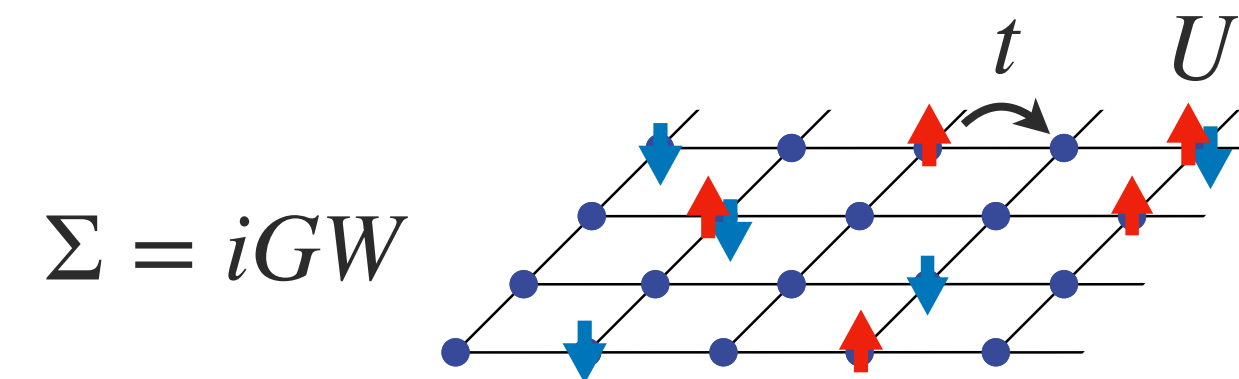
Global and causal self-consistency



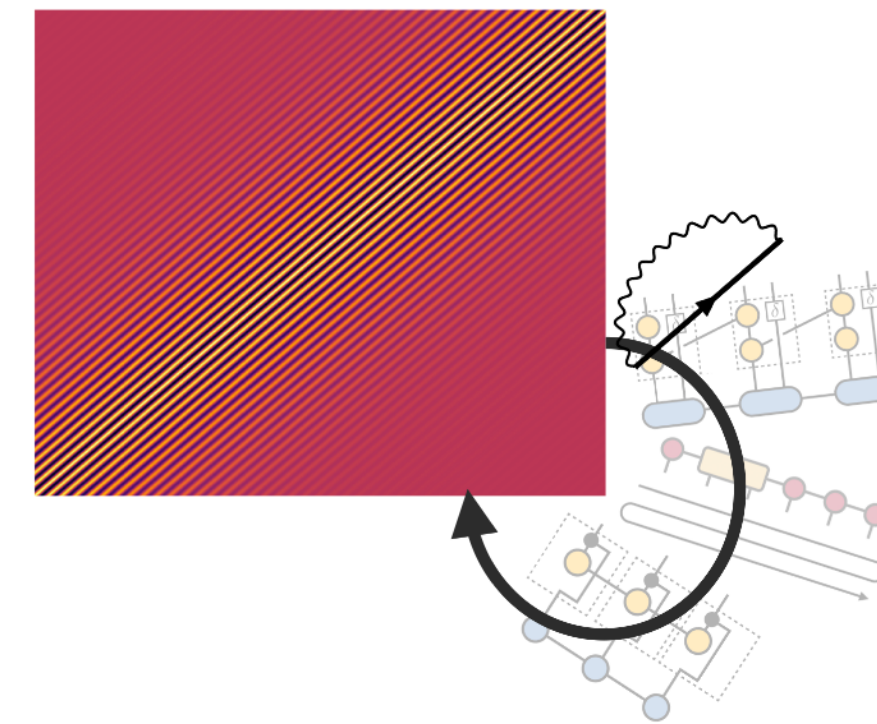
Extrapolating QTTs with dynamic mode decomposition



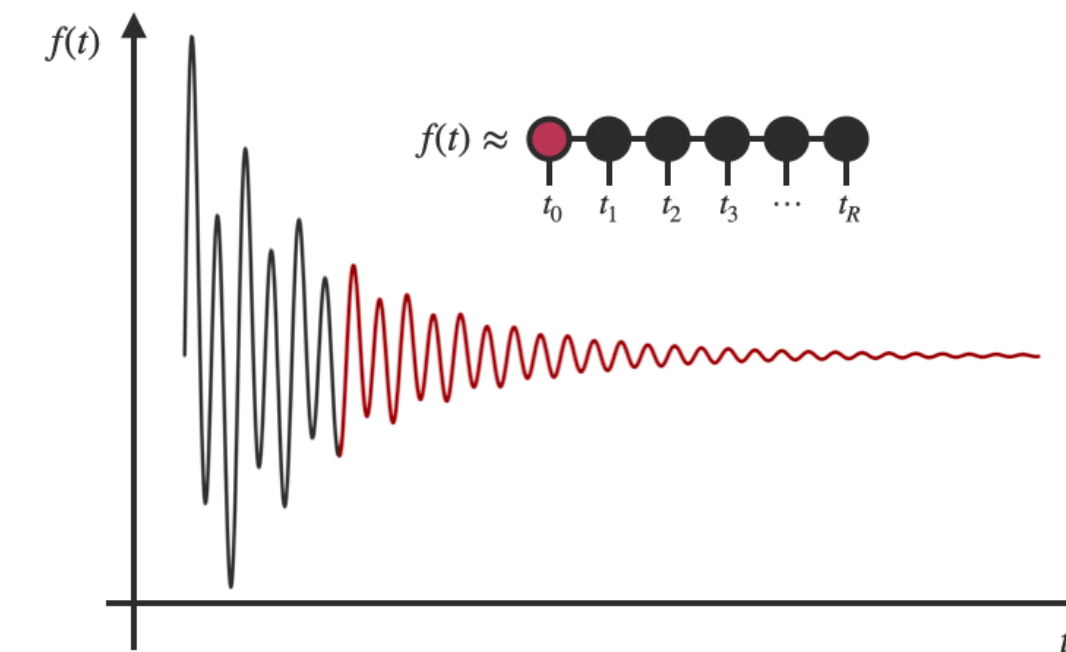
Results in application to nonequilibrium Green's functions



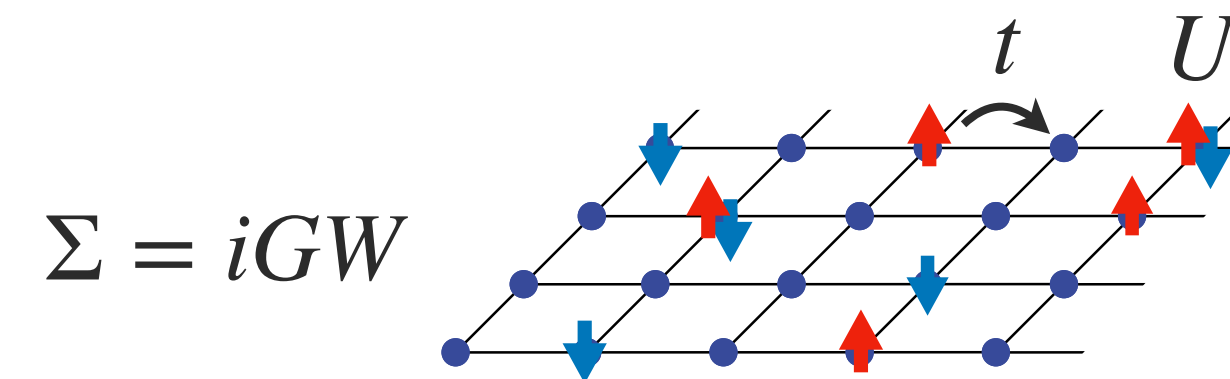
Global and causal self-consistency



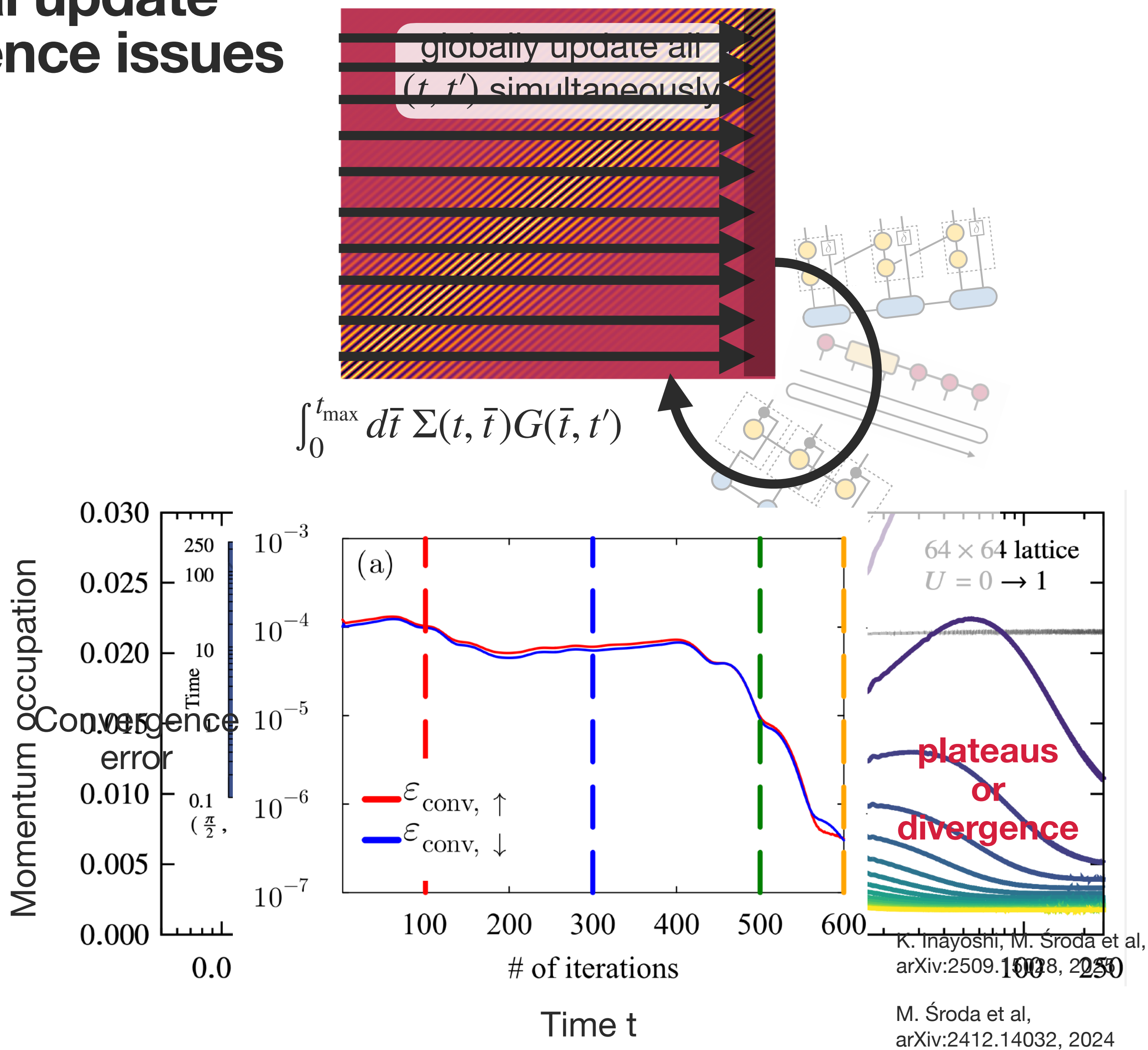
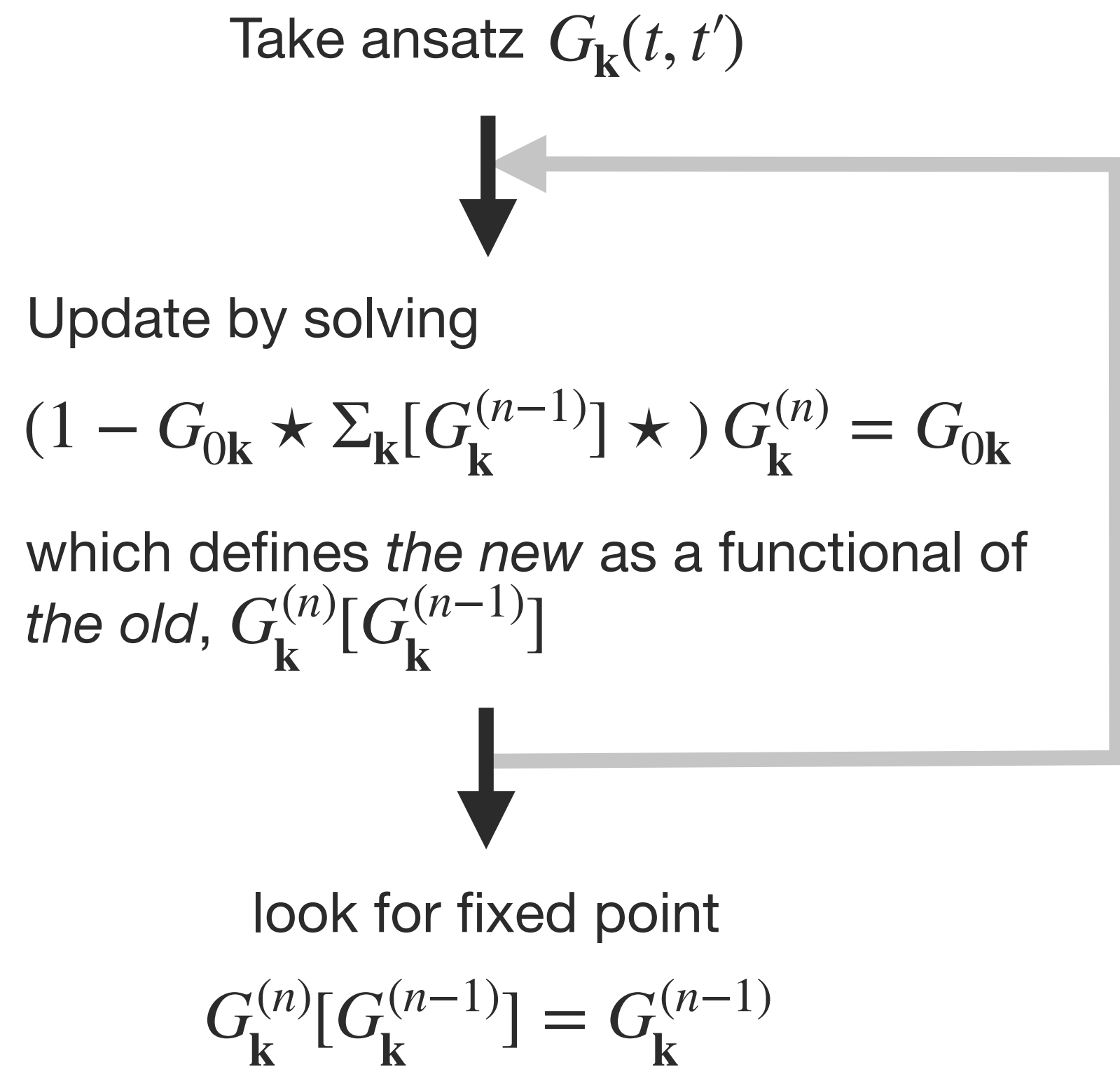
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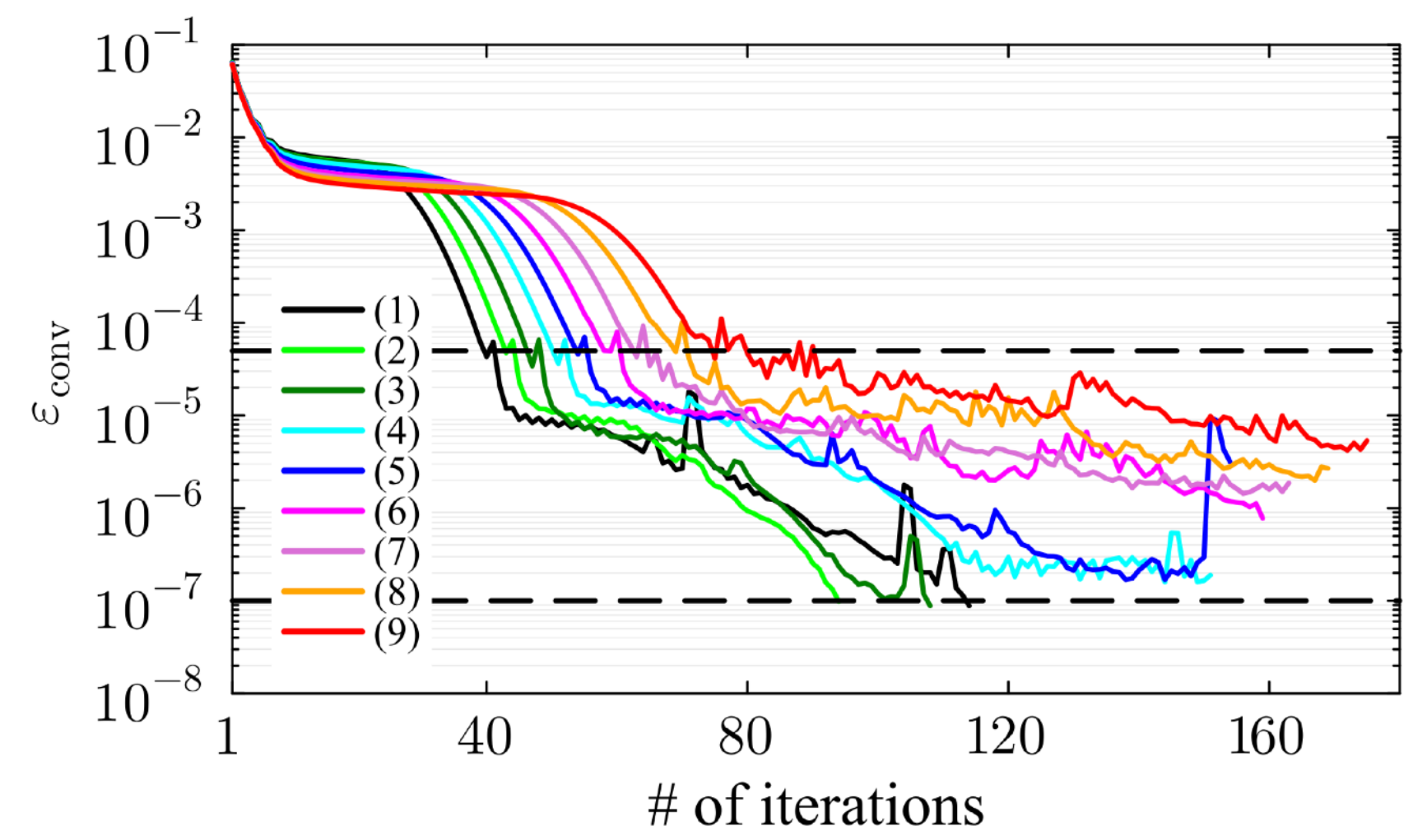
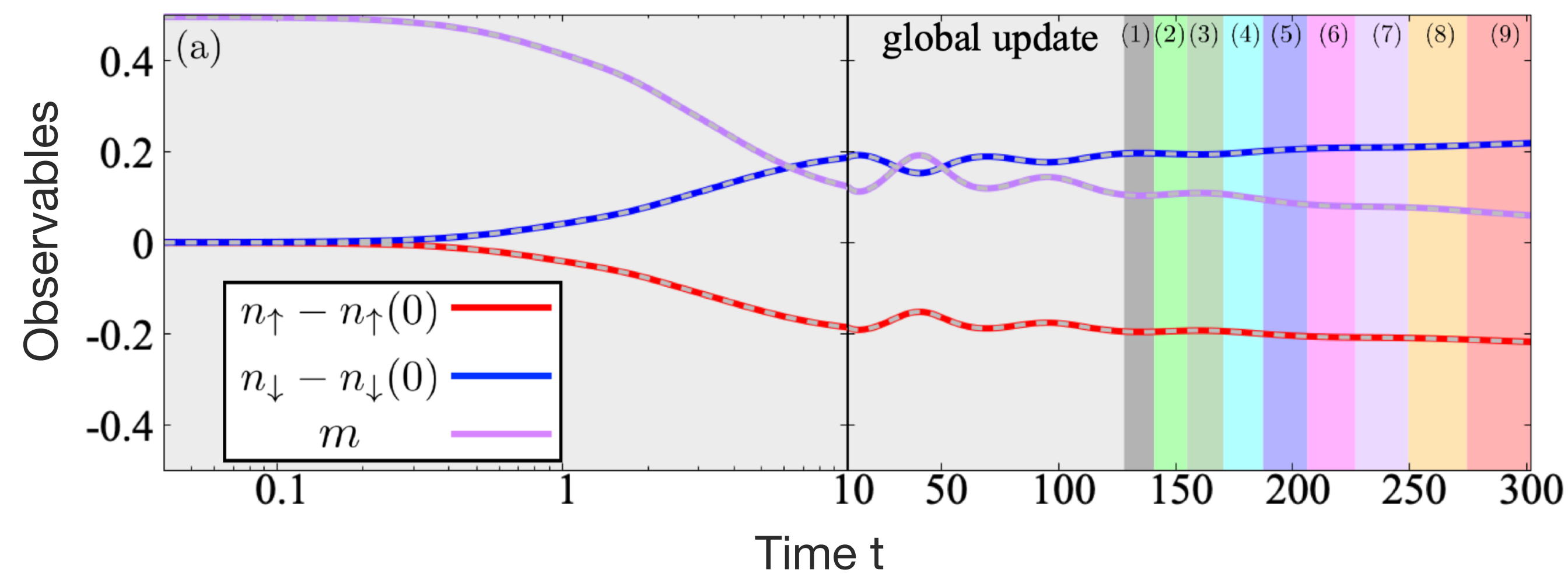
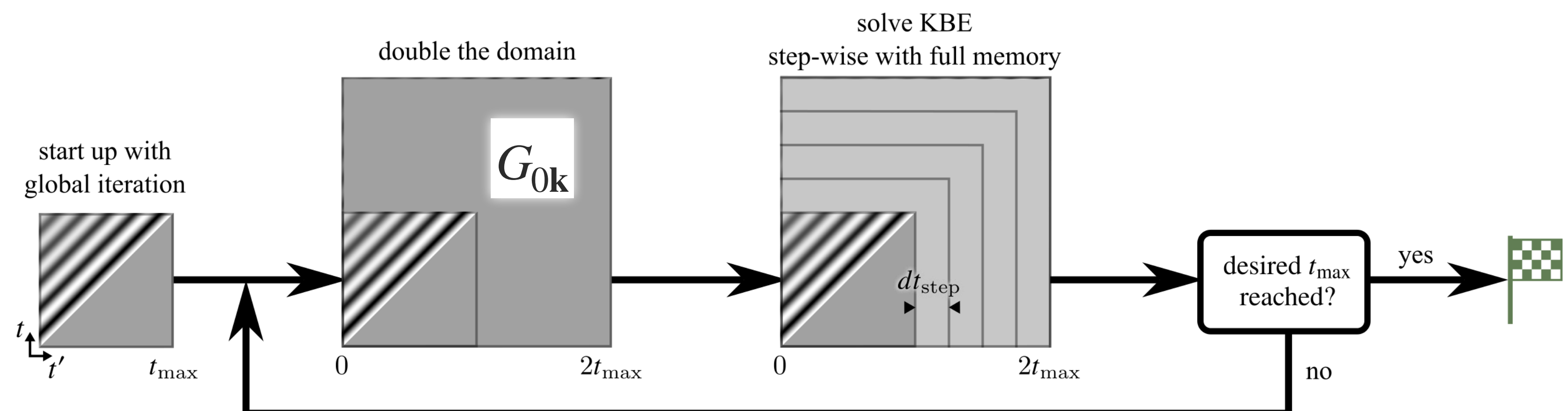
Results in application to nonequilibrium Green's functions



Simplest solver uses a global update which suffers from convergence issues

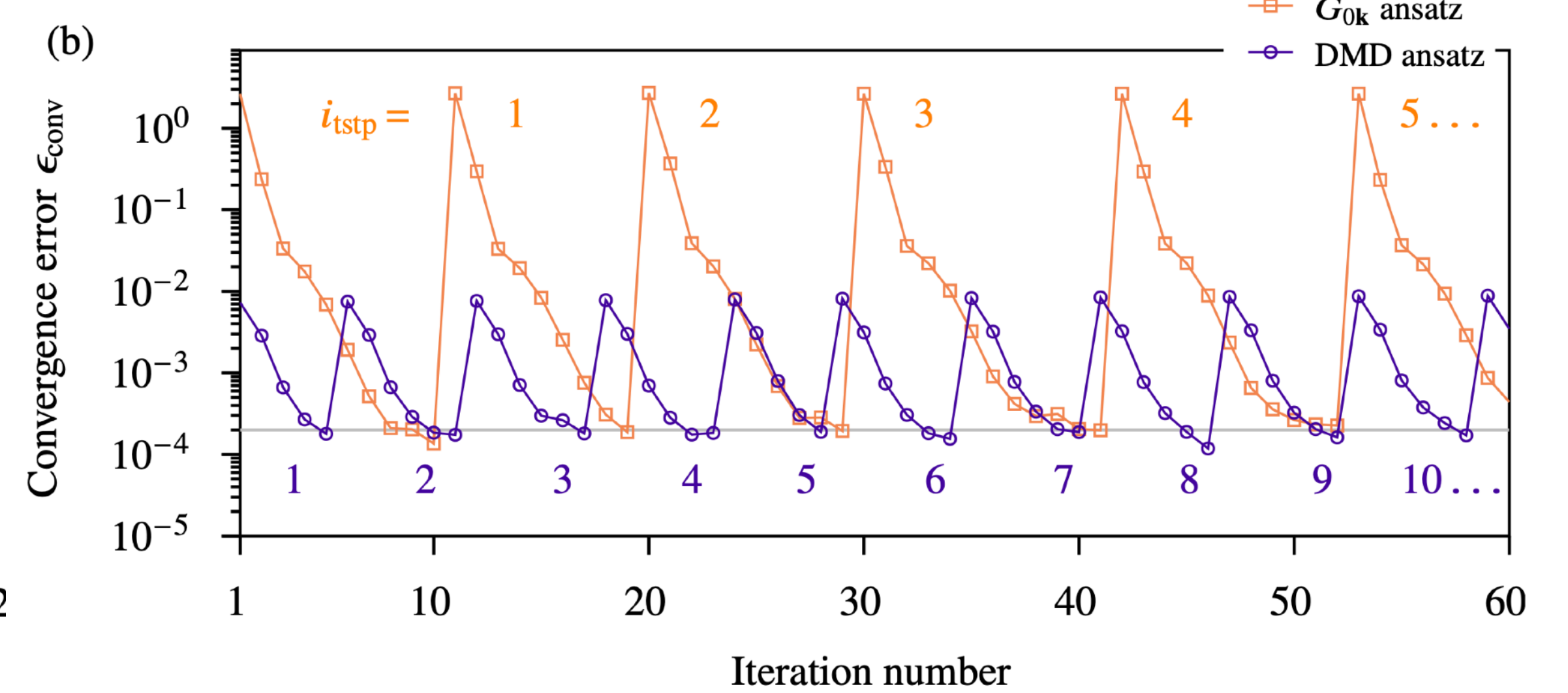
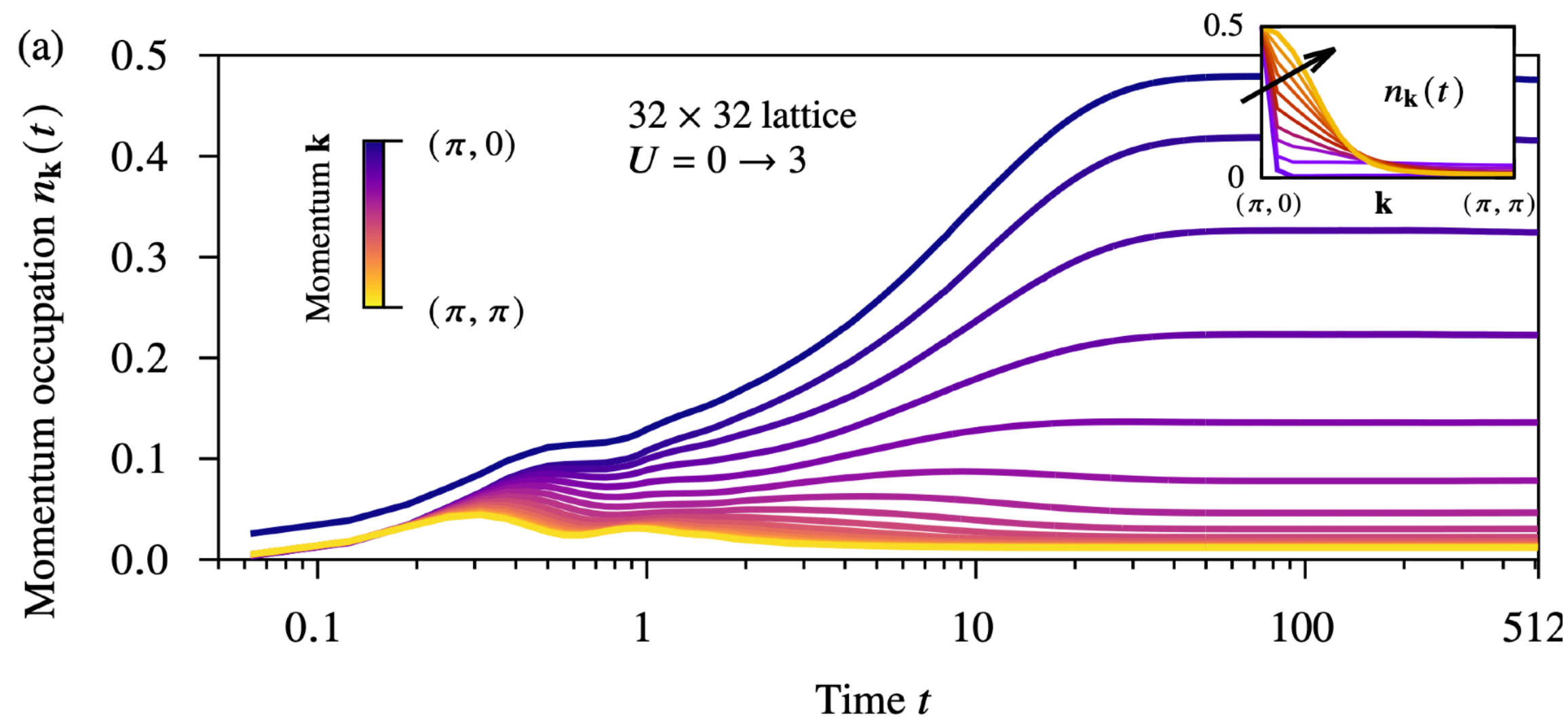
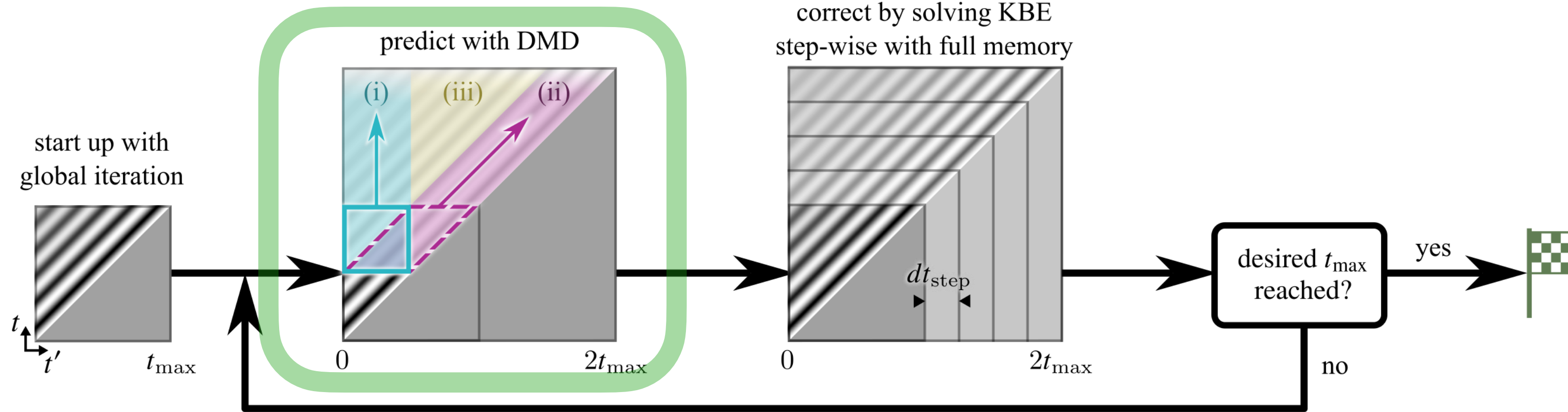


Enforcing causality makes convergence more stable



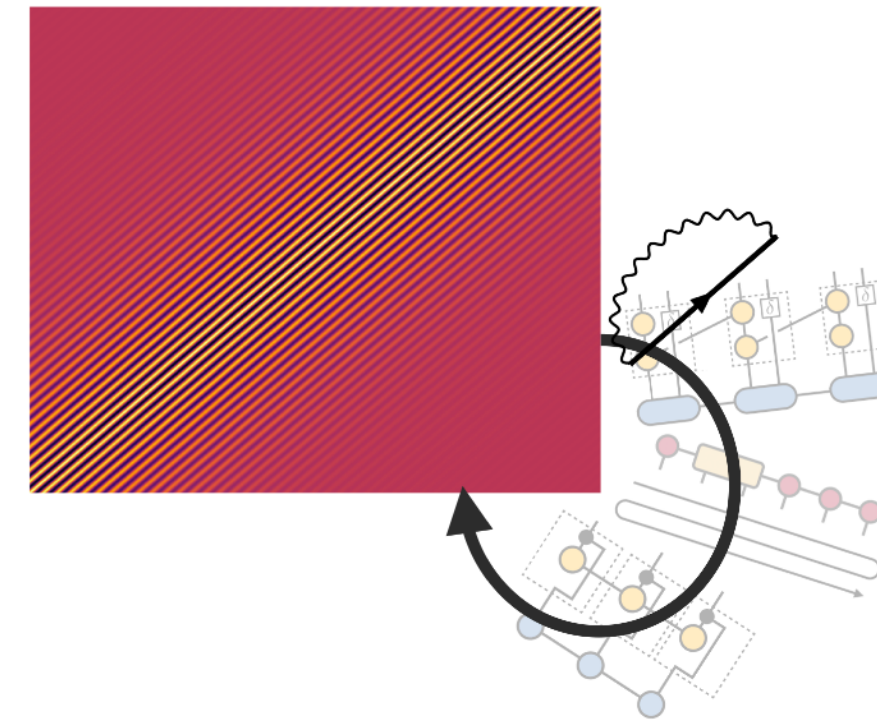
K. Inayoshi, M. Środa et al, arXiv:2509.15028, 2025

Using DMD extrapolation as initial guess gives stable, predictable and accelerated convergence

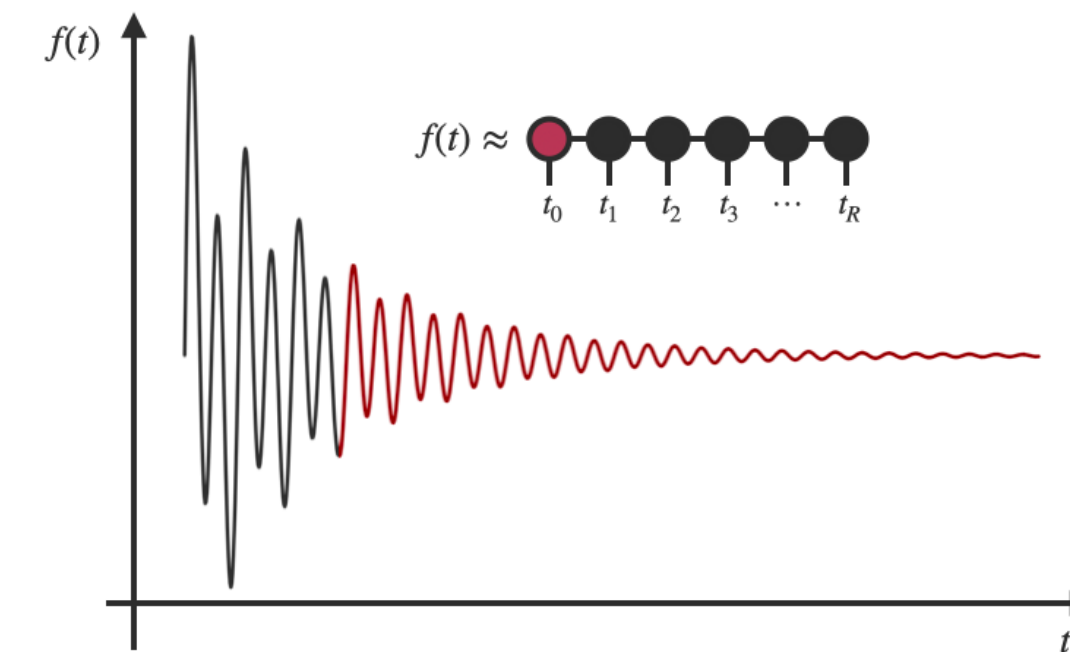


M. Środa et al, arXiv:2509.22177, 2025

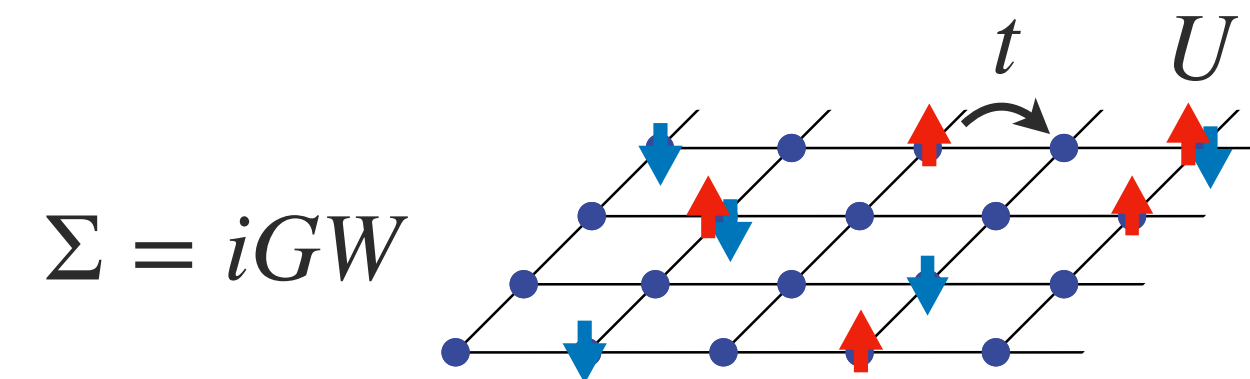
Global and causal self-consistency



Extrapolating QTTs with dynamic mode decomposition



Results in application to nonequilibrium Green's functions



Let's recall the matrix-based DMD

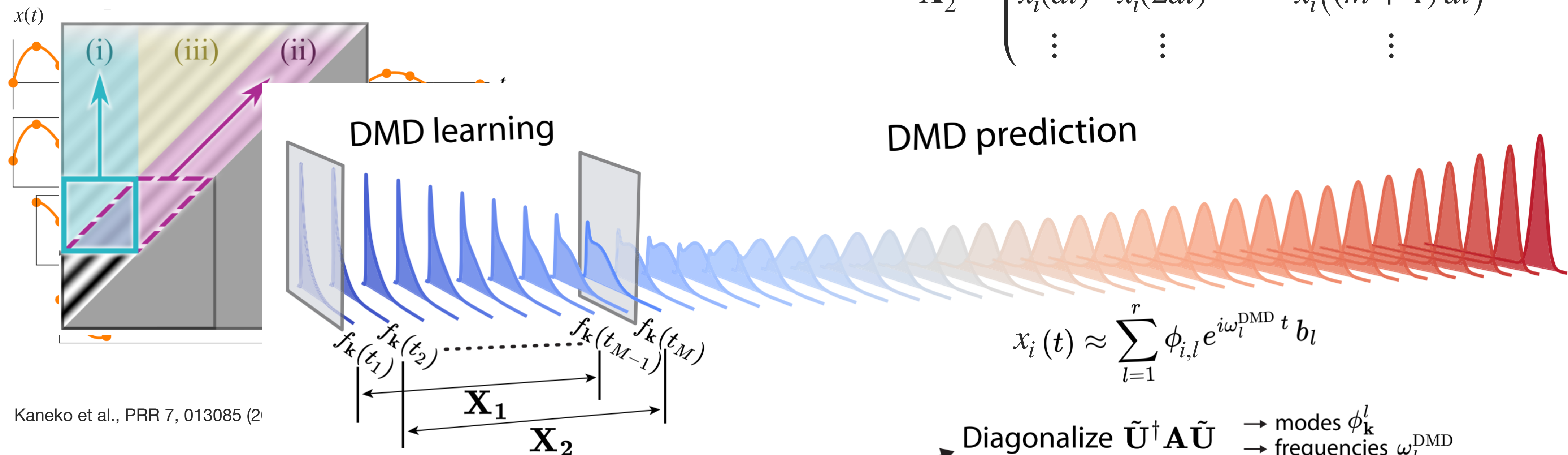
time points

observables

$$\mathbf{X} = \begin{pmatrix} \vdots & \vdots & \vdots & \vdots \\ x_i(0) & x_i(dt) & \cdots & x_i(m\,dt) & \cdots & x_i(t_{\max}) \\ \vdots & \vdots & \vdots & \vdots \end{pmatrix}$$

$$\mathbf{X}_1 = \begin{pmatrix} \vdots & \vdots & \vdots & \vdots \\ x_i(0) & x_i(dt) & \cdots & x_i(m\,dt) & \cdots & x_i(t_{\max} - dt) \\ \vdots & \vdots & \vdots & \vdots \end{pmatrix}$$

$$\mathbf{X}_2 = \begin{pmatrix} \vdots & \vdots & \vdots & \vdots \\ x_i(dt) & x_i(2dt) & \cdots & x_i((m+1)dt) & \cdots & x_i(t_{\max}) \\ \vdots & \vdots & \vdots & \vdots \end{pmatrix}$$



Kaneko et al., PRR 7, 013085 (2021)

$$\frac{d\mathbf{x}(t)}{dt} = \mathbf{A}\mathbf{x}(t) \rightarrow \mathbf{X}_2 = \mathbf{A}\mathbf{X}_1$$

$$\mathbf{A} = \mathbf{X}_2\mathbf{X}_1^{-1}$$

$$\text{SVD: } \mathbf{X}_1 \approx \tilde{\mathbf{U}}\tilde{\Sigma}\tilde{\mathbf{V}}^\dagger$$

(full operator would be $\mathbf{A} = \mathbf{X}_2\tilde{\mathbf{V}}\tilde{\Sigma}^{-1}\tilde{\mathbf{U}}^\dagger$)

$$\text{Diagonalize } \tilde{\mathbf{U}}^\dagger \mathbf{A} \tilde{\mathbf{U}} \rightarrow \begin{matrix} \text{modes } \phi_{\mathbf{k}}^l \\ \text{frequencies } \omega_l^{\text{DMD}} \\ \text{mode amplitudes } b_l \text{ from } f_{\mathbf{k}}(t_1) \end{matrix}$$

$$x_i(t) \approx \sum_{l=1}^r \phi_{i,l} e^{i\omega_l^{\text{DMD}} t} b_l$$

I. Maliyov et al., npj Computational Materials 10, 123 (2024)

Let's translate DMD into the QTT language

$$X = \begin{pmatrix} \vdots & \vdots & & \vdots & & \vdots \\ x_i(0) & x_i(dt) & \cdots & x_i(m\,dt) & \cdots & x_i(t_{\max}) \\ \vdots & \vdots & & \vdots & & \vdots \end{pmatrix}$$

Assume **quantics** encoding $m = t/dt = [t_1, \dots, t_R]_2$

$$\begin{aligned}
X &= \begin{array}{c} \boxed{} \\ \text{\scriptsize i} \quad \text{\scriptsize m} \end{array} \approx \begin{array}{c} \boxed{} - \boxed{} - \boxed{} - \cdots - \boxed{} \\ \text{\scriptsize i} \quad \text{\scriptsize t_1} \quad \text{\scriptsize t_2} \quad \quad \text{\scriptsize t_R} \end{array} \\
&\approx \begin{array}{c} \boxed{} - \boxed{} - \cdots - \boxed{} - \boxed{} - \boxed{} - \cdots - \boxed{} \\ \text{\scriptsize i_1} \quad \text{\scriptsize i_2} \quad \quad \text{\scriptsize i_N} \quad \text{\scriptsize t_1} \quad \text{\scriptsize t_2} \quad \quad \text{\scriptsize t_R} \end{array}
\end{aligned}$$

and **assume i also factorizes** in some way.



X_1, X_2 obtained by a shift with MPO
(and zeroing of the last col)

We aim to find the operator

$$AX_1 = X_2, \quad A = X_2X_1^{-1}$$

so we need **SVD** of X_1 to calculate the pseudoinverse

$$\begin{aligned}
X_1 &= \underbrace{\square \text{---} \square \text{---} \square}_{i_n} \text{---} \underbrace{\square \text{---} \square \text{---} \square}_{t_n} \\
&= \underbrace{\triangleleft \text{---} \triangleleft}_{i_n} \text{---} \underbrace{\square \text{---} \triangleright \text{---} \triangleright}_{t_n} = \underbrace{\triangleleft \text{---} \triangleleft}_{i_n} \text{---} \underbrace{\triangleleft \text{---} \alpha \text{---} \beta \text{---} \triangleright}_{S} \text{---} \underbrace{\triangleright \text{---} \triangleright}_{t_n} \\
&= \underbrace{\triangleleft \text{---} \triangleleft \text{---} \triangleleft}_{i_n}^U \text{---} \underbrace{\alpha \text{---} \diamond \text{---} \beta}_{S} \text{---} \underbrace{\triangleright \text{---} \triangleright \text{---} \triangleright}_{t_n}^{V^\dagger}
\end{aligned}$$

The pseudoinverse thus is the usual

$$X_1^{-1} = VS^{-1}U^\dagger = \text{Diagram with 6 triangles and a central diamond labeled } \beta \text{ and } \alpha, \text{ with inputs } t_{\bar{n}} \text{ and } i_{\bar{n}}.$$

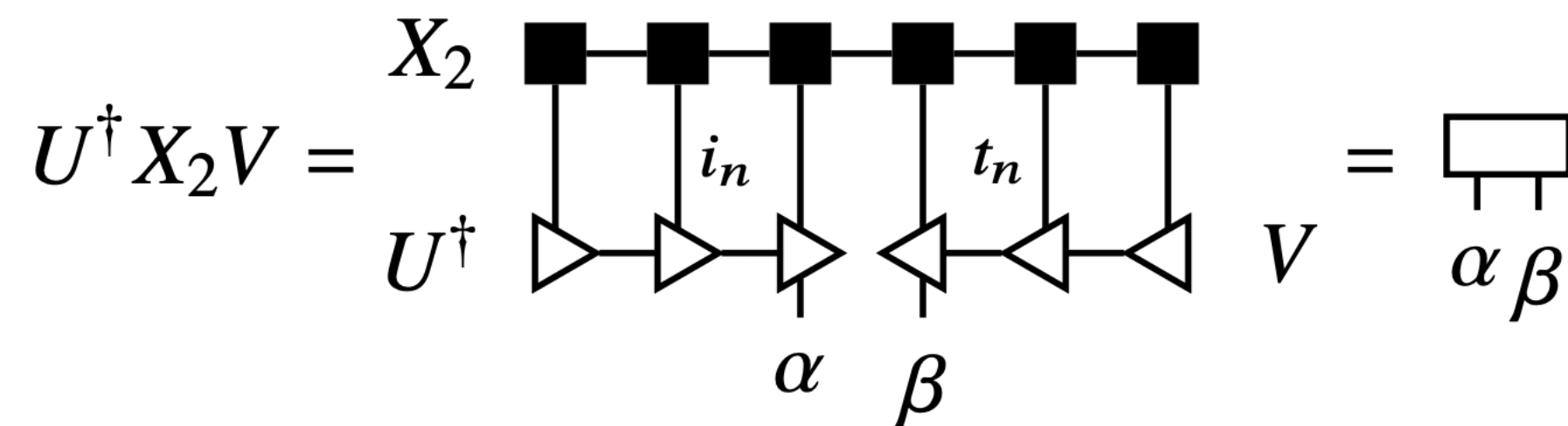
Klus et al., Nonlinearity
31, 3359–3380 (2018)

$$S^{-1} = \beta \text{---} \text{---} \alpha$$

Define **rank- r approximation** to the full operator A

$$\tilde{A} \equiv U^\dagger A U = U^\dagger X_2 V S^{-1}$$

Diagrammatically,



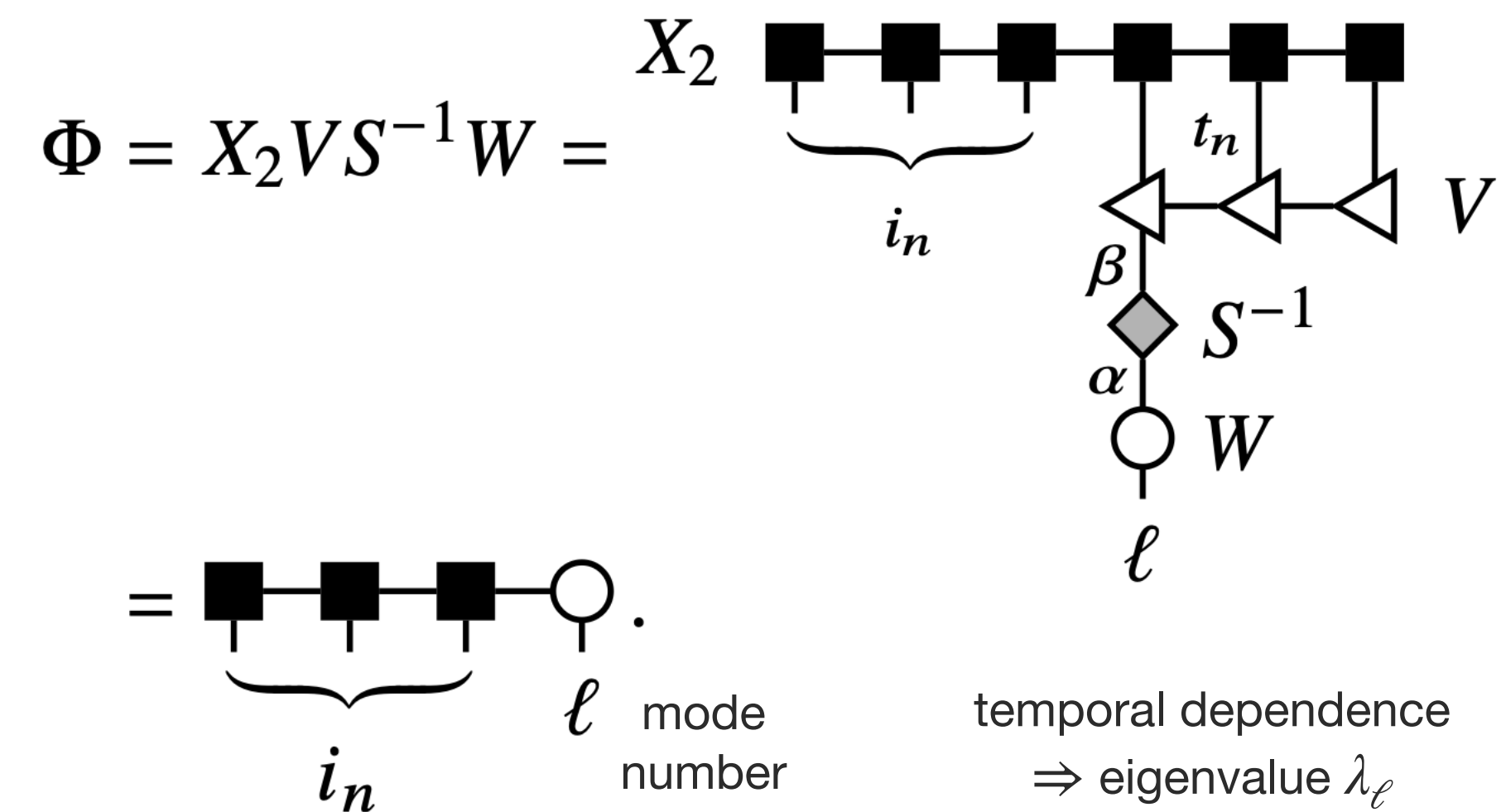
$$\tilde{A} = \alpha \text{---} \boxed{}^{\beta} \text{---} \diamond \text{---} \alpha' = \begin{array}{c} \text{---} \\ | \quad | \\ \text{---} \end{array}_{\alpha \alpha'}$$

Eigendecompose with a std dense-matrix algorithm

$$\tilde{A} = \text{bubble}(\alpha, \alpha') = \overset{W \quad \Lambda \quad W^{-1}}{\text{chain}(\alpha, \alpha')},$$

Matrices Λ , W contain approximate eigenvalues and eigenvectors in the reduced space.

The dynamic modes are the eigenvectors of full A in the original space, hence transforming back



“spatial” dependence

The **dynamic mode decomposition** is a **spatio-temporal superposition** of the modes Φ with amplitudes b ,

$$X \approx \Phi \Lambda^m b = \Phi \text{diag}(|\lambda_\ell|^m e^{i\omega_\ell m}) b,$$

$$m = t/dt = [t_1, \dots, t_R]_2$$

Amplitudes are fitted to the initial condition x

$$b = \Phi^{-1} x, \quad \Phi = \underbrace{\begin{array}{c} \blacksquare \quad \blacksquare \quad \blacksquare \\ | \quad | \quad | \\ \hline \end{array}}_{i_n} \circ \ell$$

$$b = \Phi^{-1} \begin{array}{c} \circ \\ | \\ \ell' \end{array} \begin{array}{c} \blacksquare \quad \blacksquare \quad \blacksquare \\ | \quad | \quad | \\ \square \quad \square \quad \square \end{array} x = \begin{array}{c} \text{hexagon} \\ | \\ \ell' \end{array}$$

Finally, we prepare Λ^m as a single QTT representing all m powers.

For a single eigenvalue λ , we note that

$$\lambda^m = e^{m \ln \lambda} = e^{[t_1, \dots, t_R]_2 \ln \lambda} = \prod_{n=1}^R e^{2^{R-n} t_n \ln \lambda} = \begin{array}{c} \text{pentagon}^1 \text{---} \text{pentagon}^1 \text{---} \text{pentagon}^1 \\ | \quad | \quad | \\ \ell \quad t_n \quad \ell \end{array},$$

We repeat the above for each λ , tag each QTT with a dummy tensor, and sum

$$\Lambda^m = \mathbb{I}_{\ell \Rightarrow 1} \begin{array}{c} \bullet \\ | \\ \ell \end{array} \begin{array}{c} \text{pentagon}^1 \text{---} \text{pentagon}^1 \text{---} \text{pentagon}^1 \text{---} \text{pentagon}^1 \\ | \quad | \quad | \quad | \\ \ell \quad t_n \quad \ell \end{array} + \mathbb{I}_{\ell \Rightarrow 2} \begin{array}{c} \bullet \\ | \\ \ell \end{array} \begin{array}{c} \text{pentagon}^1 \text{---} \text{pentagon}^1 \text{---} \text{pentagon}^1 \text{---} \text{pentagon}^1 \\ | \quad | \quad | \quad | \\ \ell \quad t_n \quad \ell \end{array} + \dots$$

$$= \begin{array}{c} \bullet \\ | \\ \ell \end{array} \begin{array}{c} \text{pentagon} \text{---} \text{pentagon} \text{---} \text{pentagon} \text{---} \text{pentagon} \\ | \quad | \quad | \quad | \\ \ell \quad t_n \quad \ell \end{array} \rightarrow \ell' \text{---} \begin{array}{c} \bullet \\ | \\ \ell \end{array} \begin{array}{c} \text{pentagon} \text{---} \text{pentagon} \text{---} \text{pentagon} \text{---} \text{pentagon} \\ | \quad | \quad | \quad | \\ \ell \quad t_n \quad \ell \end{array},$$

We have all the ingredients: $\Phi = \underbrace{\blacksquare \blacksquare \blacksquare}_{i_n} \circ_\ell$ $\Lambda^m = \ell' \text{---} \bullet_\ell \text{---} \text{pentagon}_{t_n} \text{---} \text{pentagon}_{t_n} \text{---} \text{pentagon}_{t_n},$ $b = \text{hexagon}_{\ell'}$

So we perform the final contraction:

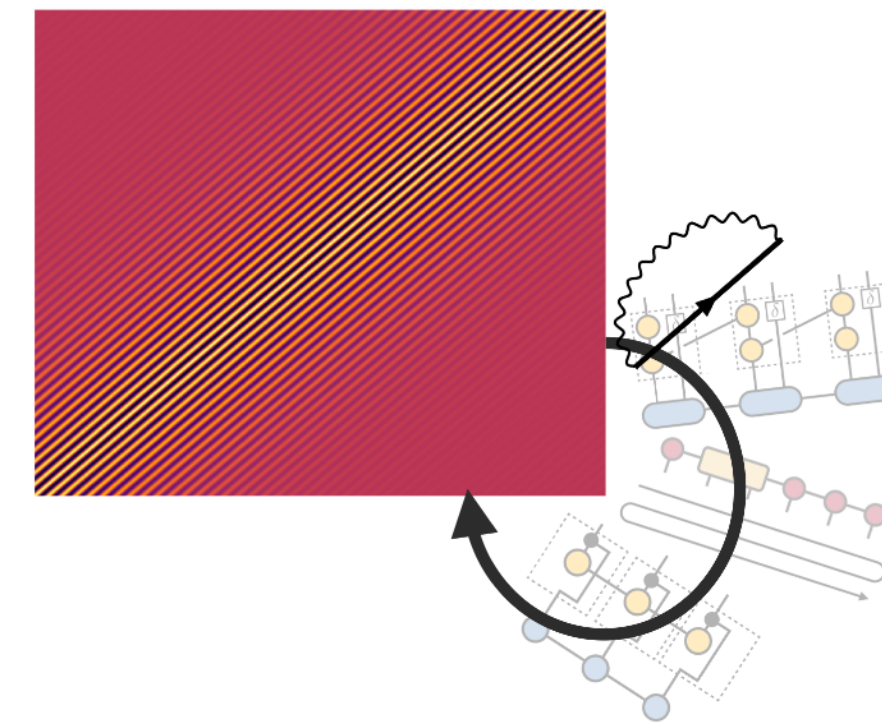
$$X \approx \Phi \Lambda^m b = \Phi \text{diag}(|\lambda_\ell|^m e^{i\omega_\ell m}) b,$$

$$X \approx \underbrace{\blacksquare \blacksquare \blacksquare}_{i_n} \text{---} \circ_\ell \text{---} \bullet_{\ell'} \text{---} \text{pentagon}_{t_n} \text{---} \text{pentagon}_{t_n} \text{---} \text{pentagon}_{t_n}.$$

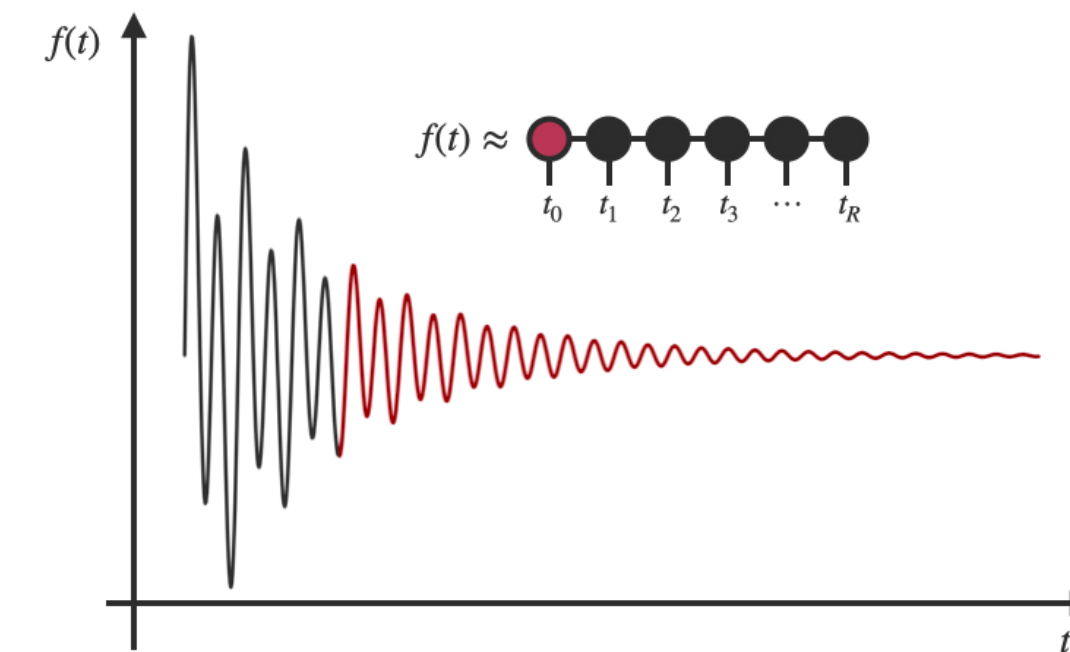
$$= \underbrace{\blacksquare \blacksquare \blacksquare}_{i_n} \text{---} \text{pentagon}_{t_n} \text{---} \text{pentagon}_{t_n} \text{---} \text{pentagon}_{t_n}.$$

- ✓ “recompression”
- ✓ denoising (set cutoff appropriately in initial SVD of X_1)
- ✓ extrapolation (add coarse tensor in Λ^m)
- ✓ interpolation (add fine tensor in Λ^m)
- ✓ DMD is general (not only time dynamics, no smoothness requirement)

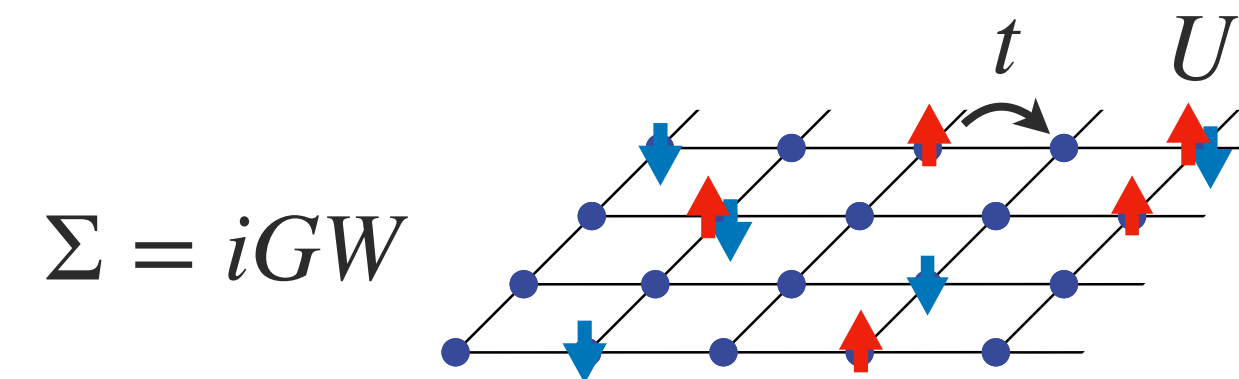
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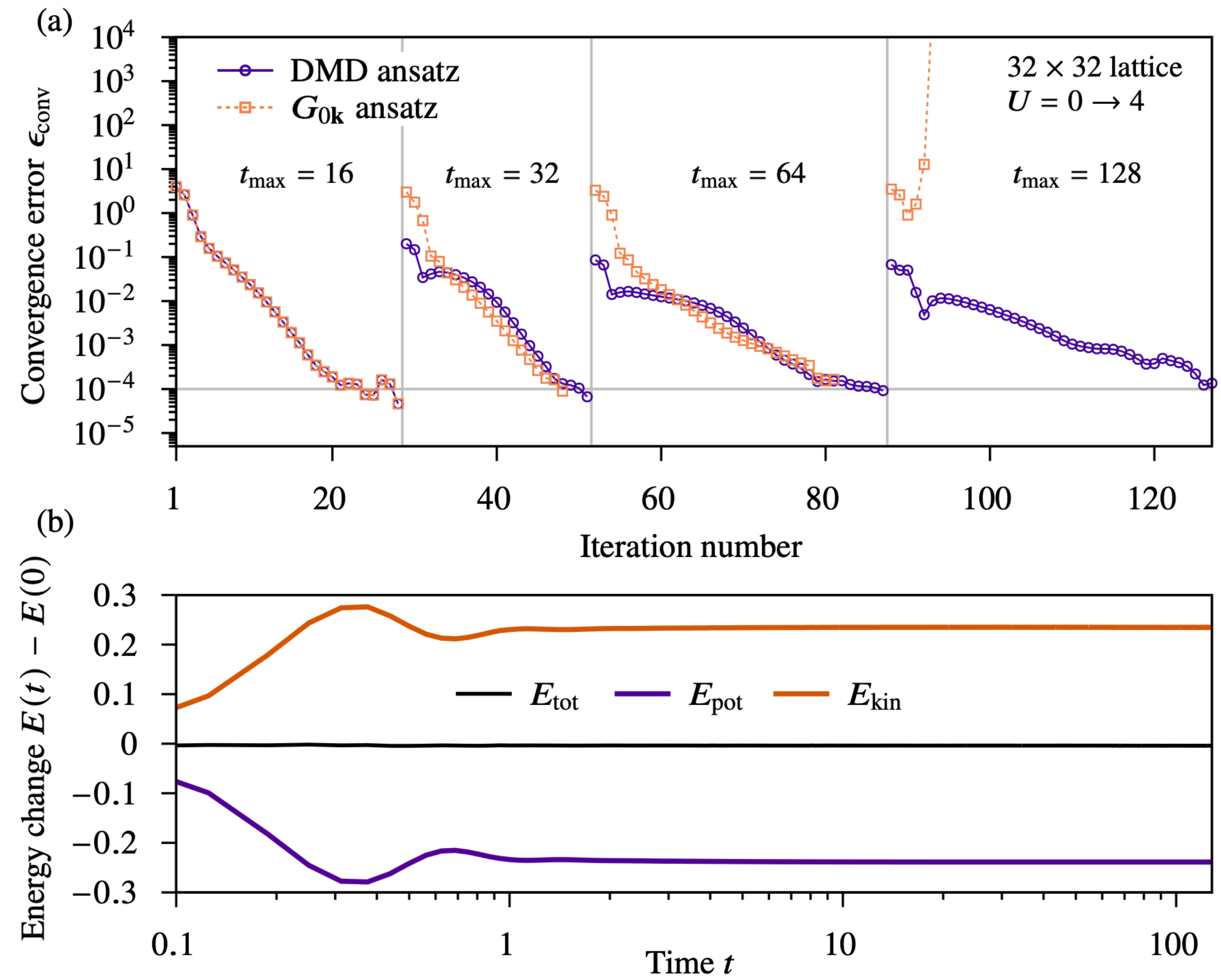
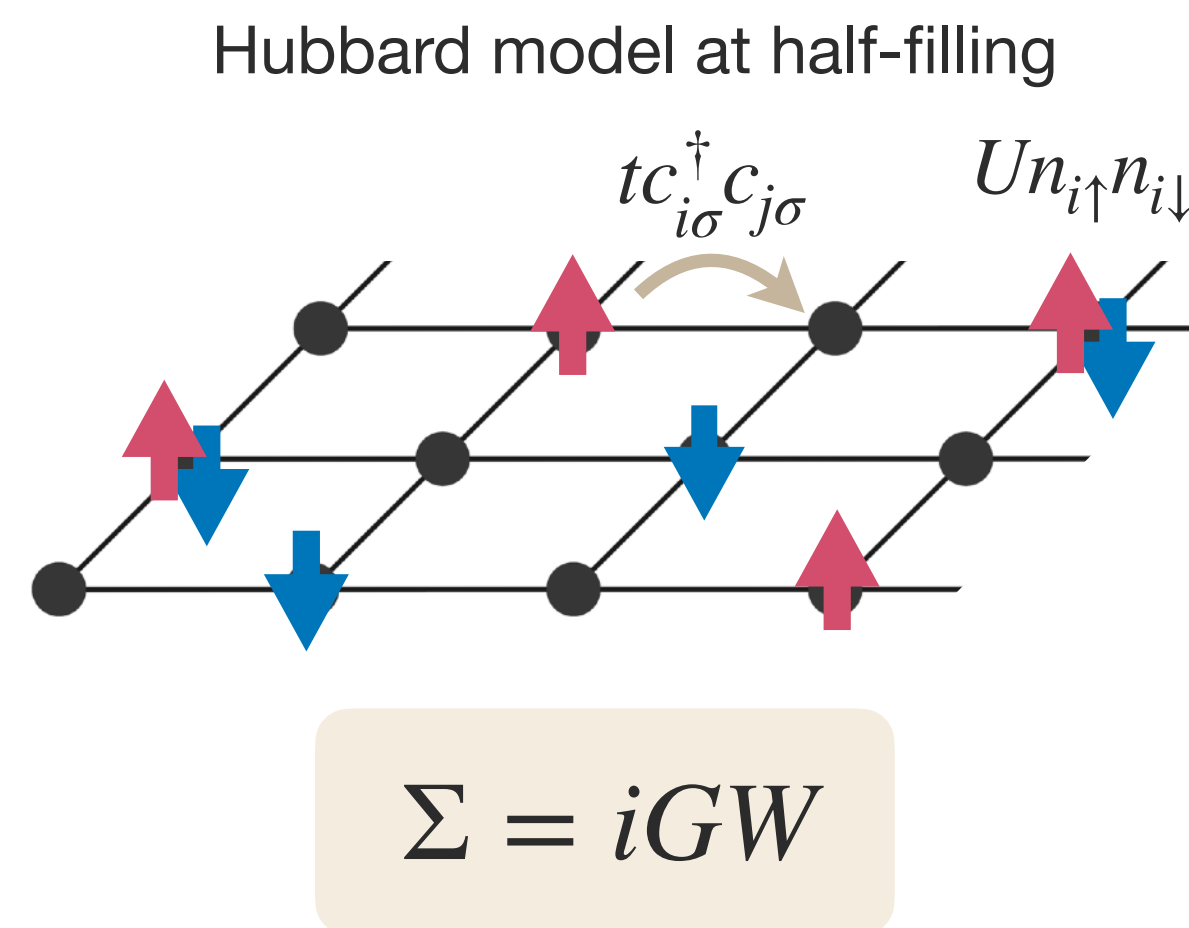
Extrapolating QTTs with dynamic mode decomposition



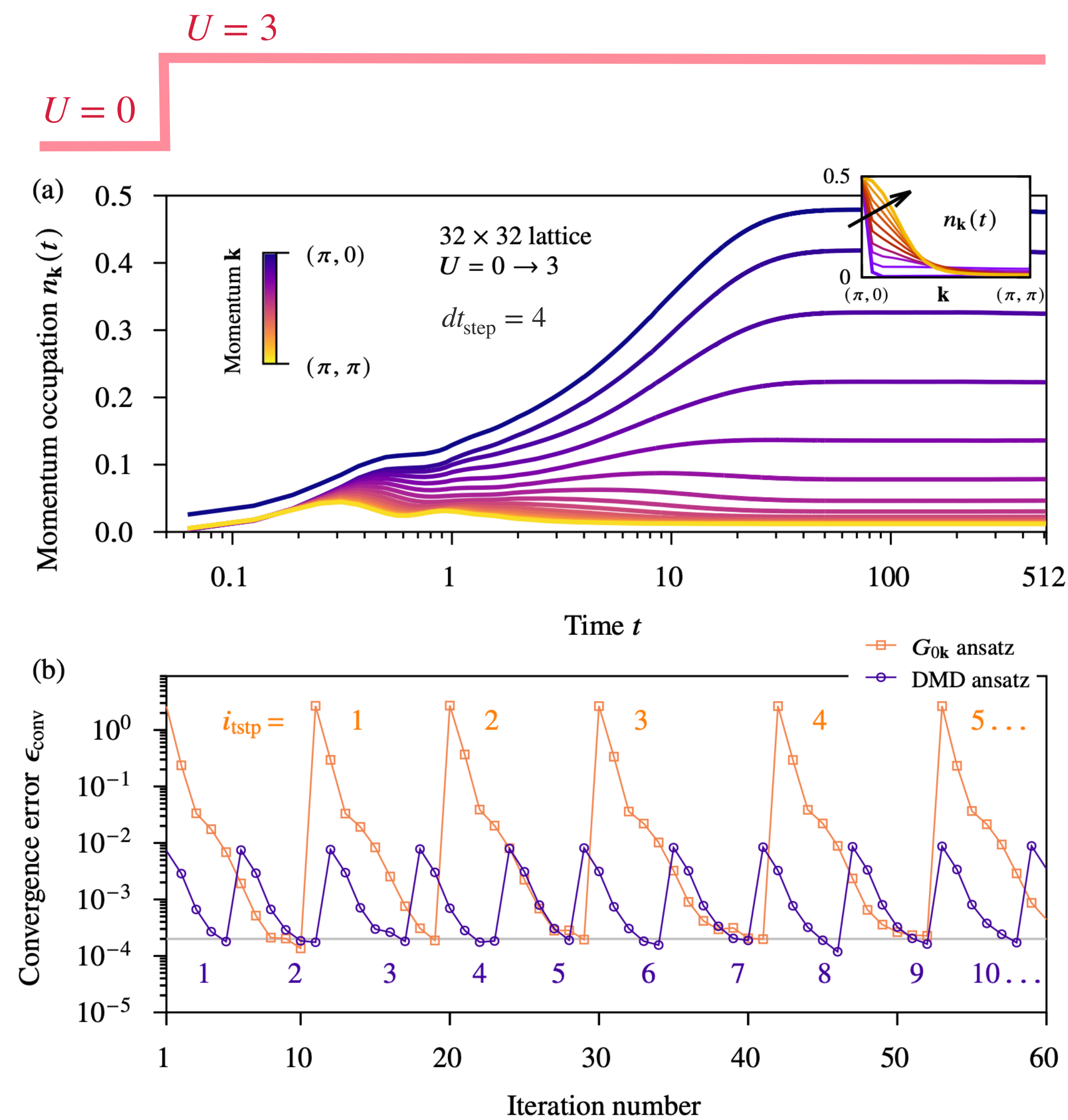
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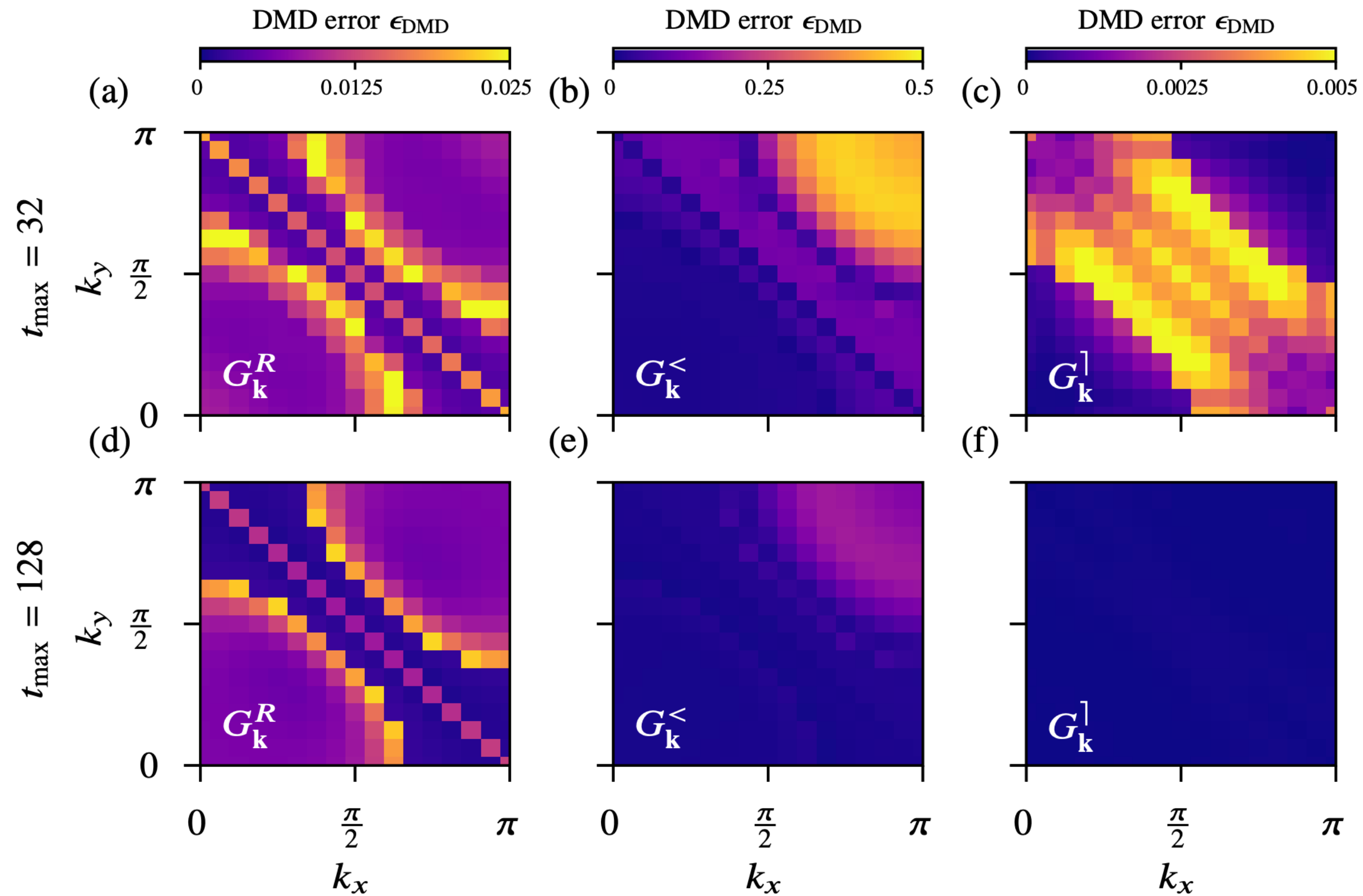
DMD guess stabilizes convergence



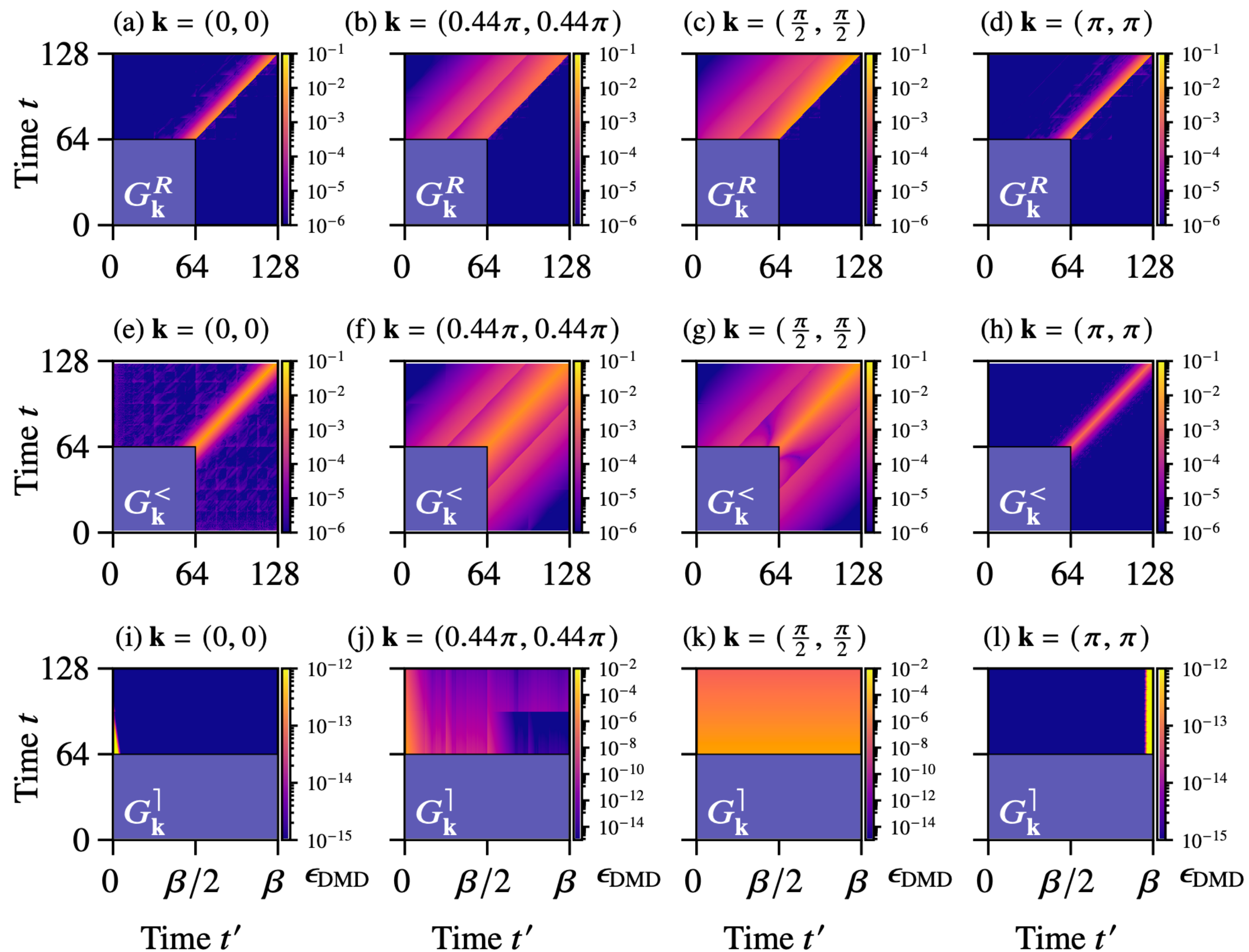
DMD guess accelerates convergence



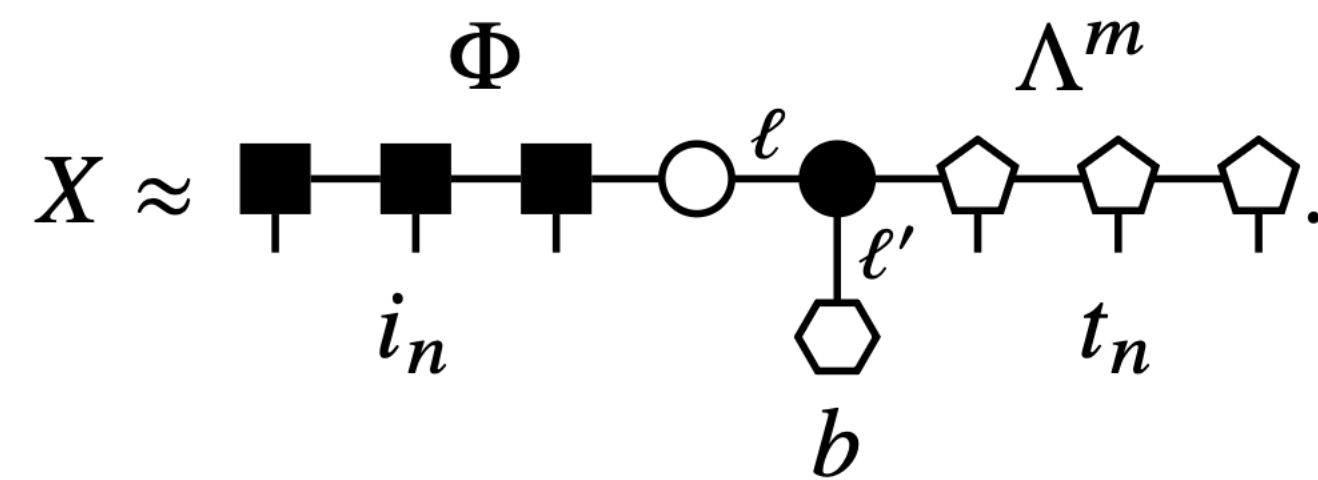
DMD guess is accurate across the BZ



DMD guess is accurate in the two-time plane



DMD is fully composable with the QTT framework and the algorithm is straightforward



DMD integrates into the QTT framework, maintaining exponentially fine resolution, and offering **extrapolation, interpolation, and denoising**

QTT-DMD provides reliable extrapolation for NEGF problems, accelerating convergence

